

CHAPTER 8. DIMENSIONAL ANALYSIS AND SIMILITUDE

1. Necessity of experimental measurements with scaled models
2. Homogeneity: Buckingham Pi theorem
3. Mechanical similitude of fluid flows

1. Necessity of experimental measurements with scaled models

- ✓ Different forms to solve a problem of Fluid Mechanics or Hydraulic Engineering:
 - Analytical: Sometimes the solutions are unknown
 - Computational: computation capacity is not infinite, necessity of empirical correlation
 - Experimental: necessity of laboratory experiments with models or mock-ups.

2. Buckingham Pi theorem

✓ Formulation

- Physical phenomenon: $f(A, B, C, \dots, N) = 0$
- "n" Physical variables: (A, B, C, \dots, N)
- "m" fundamental dimensions: usually 3 in fluid mechanics: M, L y T
- It can be defined: $F(\Pi_1, \Pi_2, \dots, \Pi_{n-m}) = 0$

With (n-m) dimensionless groups:

$$\Pi_i = A^{\alpha_i} B^{\beta_i} C^{\gamma_i} \dots N^{\nu_i}$$

2. Buckingham Pi theorem

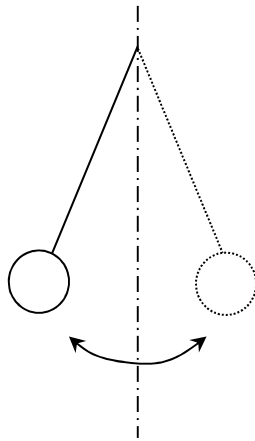
- ✓ Method of repeating variables
 1. Physical variables A, B, C, \dots, N
(m) fundamental dimensions
($n-m$) number of dimensionless groups
 2. Selection of repeating variables (m)
 3. Generate the dimensionless groups Π
 4. Calculation of the exponents

2. Buckingham Pi theorem

- ✓ Useful guidelines
 - Dimensionless physical variable, dimensionless Π group
 - Variables with the same dimensions, quotient dimensionless group
 - Any Π number can be substituted by:
 - a power of itself
 - the product by a constant
 - the quotient of two Π groups
 - Any Π number can be expressed as a function of other Π numbers

2. Buckingham Pi theorem

EXAMPLE 1: SIMPLE PENDULUM



STEP 1:
 $n = 4$ variables
 $t, [T]$
 $l, [L]$
 $m, [M]$
 $W, [MLT^{-2}]$

$m=3$ fundamental dimensions: M, L, T

N° of dimensionless groups: $n-m = 4-3 = 1$

STEP 2:
 $m = 3$ repeating variables are selected: t, l, m

STEP 3:
 Dimensionless group: $\Pi = t^{\alpha} l^{\beta} m^{\gamma} W$

STEP 4:

$$[\Pi] = [T]^{\alpha} [L]^{\beta} [M]^{\gamma} [MLT^{-2}]^1$$

$$\Pi = t^2 l^{-1} m^{-1} W$$

$$t = \left(\frac{\Pi l m}{W} \right)^{1/2} = \Pi^{1/2} \left(\frac{l}{g} \right)^{1/2}$$

$$t = 2\pi \sqrt{\frac{l}{g}}$$

2. Buckingham Pi theorem

EXAMPLE 2: DRAG RESISTANCE THAT A FLUID EXERTS ON THE MOTION OF A PARTIALLY SUBMERGED BODY

STEP 1: Acting variables: $n = 7$ variables

$$F [MLT^{-2}] \quad U [LT^{-1}] \quad \varepsilon [L] \quad g [LT^{-2}] \quad l [L] \quad \rho [ML^{-3}] \quad \mu [ML^{-1}T^{-1}]$$

$m=3$ fundamental dimensions: T, L y M

No of dimensionless groups: $n-m = 7-3 = 4$

STEP 2: Repeating variables: ρ, l, U

STEP 3: $(n-m) = 4$ Π groups.

$$\Pi_1 = \rho^{\alpha_1} l^{\beta_1} U^{\gamma_1} F$$

$$\Pi_2 = \rho^{\alpha_2} l^{\beta_2} U^{\gamma_2} \varepsilon$$

$$\Pi_3 = \rho^{\alpha_3} l^{\beta_3} U^{\gamma_3} \mu$$

$$\Pi_4 = \rho^{\alpha_4} l^{\beta_4} U^{\gamma_4} g$$

$$\text{STEP 4: } [\Pi_1] = [M L^{-3}]^{\alpha_1} [L]^{\beta_1} [LT^{-1}]^{\gamma_1} [MLT^{-2}]$$

$$[\Pi_2] = [M L^{-3}]^{\alpha_2} [L]^{\beta_2} [LT^{-1}]^{\gamma_2} [L]$$

$$[\Pi_3] = [M L^{-3}]^{\alpha_3} [L]^{\beta_3} [LT^{-1}]^{\gamma_3} [ML^{-1}T^{-1}]$$

$$[\Pi_4] = [M L^{-3}]^{\alpha_4} [L]^{\beta_4} [LT^{-1}]^{\gamma_4} [LT^{-2}]$$



$$\Pi_1 = \frac{F}{\rho l^2 U^2} \quad \Pi_2 = \frac{\varepsilon}{l}$$

$$\Pi_3' = \Pi_3^{-1} = \frac{\rho l U}{\mu} = Re$$

$$\Pi_4' = \Pi_4^{-1/2} = \frac{U}{\sqrt{lg}} = Fr$$

2. Buckingham Pi theorem

EXAMPLE 2: DRAG RESISTANCE THAT A FLUID EXERTS ON THE MOTION OF A PARTIALLY SUBMERGED BODY

- Any Π number can be expressed as a function of other Π numbers

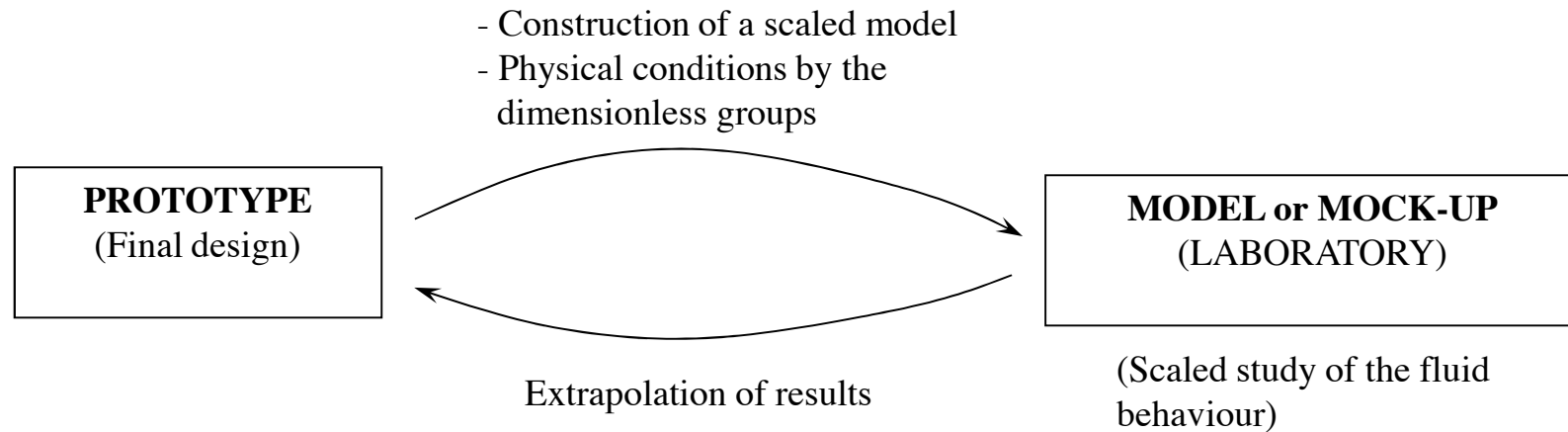
$$\boxed{\Pi_1 = \varphi(\Pi_2, \Pi_3, \Pi_4)} \quad \boxed{\frac{F}{\rho l^2 U^2} = \varphi\left(\frac{\varepsilon}{l}, \text{Re}, \text{Fr}\right)} \Rightarrow \boxed{F = \rho l^2 U^2 \cdot \varphi\left(\frac{\varepsilon}{l}, \text{Re}, \text{Fr}\right)}$$

$$\Rightarrow \boxed{F = \frac{1}{2} C_d \rho A U^2}$$

$$\Rightarrow \boxed{F = F_f + F_w = \frac{1}{2} C_{d,f} \rho A U^2 + \frac{1}{2} C_{d,w} \rho A U^2}$$

3. Mechanical similitude of fluid flows

- ✓ Study of the fluid related phenomena in the laboratory



- ✓ Mechanical or absolute similitude

- Geometric similitude
- Kinematic similitude
- Dynamic similitude

- ✓ Implications for the laboratory experiments:

$$\Pi_{1m} = \varphi (\Pi_{2m}, \Pi_{3m}, \dots, \Pi_{km}) = \varphi (\Pi_{2p}, \Pi_{3p}, \dots, \Pi_{kp}) = \Pi_{1p}$$

3.1. Dimensionless parameters in Fluid Mechanics

✓ Calculation of the main dimensionless groups

- Equations

$$\rho \vec{F} - \vec{\nabla} p + \frac{1}{3} \mu \vec{\nabla}(\vec{\nabla} \cdot \vec{U}) + \mu \nabla^2 \vec{U} = \rho \vec{a} \quad \text{NAVIER - STOKES EQUATION}$$

$$p = C \rho^k \quad \text{PROCESS EQUATION}$$

$$\frac{\partial \rho}{\partial t} + \vec{\nabla}(\rho \vec{U}) = 0 \quad \text{CONTINUITY EQUATION}$$

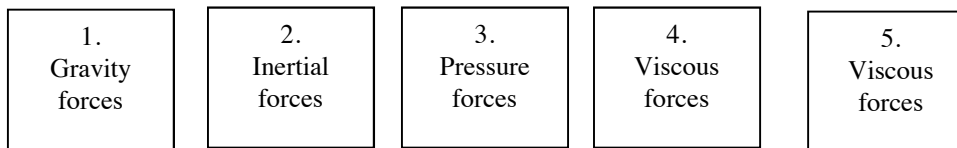
The same scale relationship between terms of the same equation in PROTOTYPE and MODEL

3.1. Dimensionless parameters in Fluid Mechanics

- ✓ Calculation of the main dimensionless groups

NAVIER – STOKES EQUATION

$$\frac{\rho_p \mathbf{g}_p}{\rho_m \mathbf{g}_m} = \frac{\rho_p \mathbf{a}_p}{\rho_m \mathbf{a}_m} = \frac{\nabla p_p}{\nabla p_m} = \frac{\mu_p \nabla(\nabla U_p)}{\mu_m \nabla(\nabla U_m)} = \frac{\mu_p \nabla^2 U_p}{\mu_m \nabla^2 U_m}$$



1. Gravity forces	2. Inertial forces	➔	$\frac{U_m}{\sqrt{g_m l_m}} = \frac{U_p}{\sqrt{g_p l_p}} = Fr$	FROUDE
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2. Inertial forces	3. Pressure forces	➔	$\frac{p_m}{\rho_m U_m^2} = \frac{p_p}{\rho_p U_p^2} = Eu$	EULER
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2. Inertial forces	4. Viscous forces	➔	$\frac{\rho_p U_p l_p}{\mu_p} = \frac{\rho_m U_m l_m}{\mu_m} = Re$	REYNOLDS
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3.1. Dimensionless parameters in Fluid Mechanics

- ✓ Calculation of the main dimensionless groups

PROCESS EQUATION

$$\begin{array}{ccc}
 \left. \begin{array}{l} \frac{p_{1p}}{p_{2p}} = \left(\frac{\rho_{1p}}{\rho_{2p}} \right)^{k_p} \\ \frac{p_{1m}}{p_{2m}} = \left(\frac{\rho_{1m}}{\rho_{2m}} \right)^{k_m} \end{array} \right\} & & \\
 \left. \begin{array}{l} \frac{p_{1m}}{p_{2m}} = \frac{p_{1p}}{p_{2p}} \\ \frac{\rho_{1m}}{\rho_{2m}} = \frac{\rho_{1p}}{\rho_{2p}} \end{array} \right\} & \Rightarrow & \boxed{\frac{k_m}{k_p} = 1} \quad \text{"STATE CONDITION"}
 \end{array}$$

3.1. Dimensionless parameters in Fluid Mechanics

- ✓ Calculation of the main dimensionless groups

PROCESS EQUATION

$$c^2 = \frac{dp}{d\rho} = \frac{kp}{\rho}$$

NEWTON'S VELOCITY LAW

$$\frac{c_m^2}{c_p^2} = \frac{k_m p_m / \rho_m}{k_p p_p / \rho_p} \quad \Rightarrow \quad \frac{c_m^2 U_p^2}{c_p^2 U_m^2} = \frac{k_m}{k_p} \frac{p_m}{\rho_m U_m^2} \frac{\rho_p U_p^2}{p_p} = \frac{Eu_m}{Eu_p} = 1$$

$$\Rightarrow \frac{c_m^2 U_p^2}{c_p^2 U_m^2} = 1 \quad \Rightarrow \quad \boxed{\frac{U_p}{c_p} = \frac{U_m}{c_m} = M} \quad \text{MACH}$$

3.2. Physical meaning of dimensionless groups

NAME	EQUATION	MEANING	RELEVANCE
Reynolds	$Re = \frac{Ul}{\nu} = \frac{\rho Ul}{\mu}$	Relationship between inertial forces and viscous forces	- Differentiating the flow regime (turbulent / laminar). - Viscous effects
Froude	$Fr = \frac{U}{\sqrt{gl}}$	Relationship between inertial forces and gravity forces	- Open-channel flows
Mach	$M = \frac{U}{c}$	Relationship between inertial forces and elastic forces	- Compressible flow - Sound propagation (pressure waves)
Euler (Cavitation No if $p-p_v$ instead of p)	$Eu = \frac{p}{\rho U^2}$	Relationship between pressure forces and inertial forces	- Cavitation phenomenon
Thermodynamic process exponent	k	Type of thermodynamic process a gas develops	- Compressible flow (gases)
Weber	$We = \frac{U^2 \rho L}{\sigma}$	Relationship between inertial forces and surface tension forces	- Flow with an interface between two liquids
Strouhal	$S = \frac{\omega l}{U}$	Importance of undulatory motion (angular velocity ω) in comparison to the main motion (average vel. U)	- Oscillating flows
Rossby	$Ro = \frac{U}{\Omega_{tierra} l}$	Relative importance of the Coriolis effect in the fluid flow	- geophysical flows
Relative roughness	$\frac{\epsilon}{l}$	Roughness with respect to the main length of the body	- Turbulent flow - Friction
Friction coefficient or Fanning friction factor	$C_f = \frac{\tau_p}{\frac{1}{2}\rho U^2}$	Relationship between shear friction stress (per unit area) and dynamic pressure	- Friction
Friction factor	$f = \frac{h_f}{(U^2 / 2g)(L/D)}$	Relationship between head loss and kinetic head	- Internal viscous flow (pipes)
Drag coefficient	$C_d = \frac{F_d}{(1/2)\rho AU^2}$	Relationship between drag force and dynamic force	- External flow (aerodynamics, hydrodynamics)
Lift coefficient	$C_L = \frac{F_L}{(1/2)\rho AU^2}$	Relationship between lift force and dynamic force	- External flow (aerodynamics, hydrodynamics)

3.3. Absolute and limited similitude. Froude and Reynolds incompatibility

✓ Mechanical or absolute similitude

- Geometric similitude
- Kinematic similitude
- Dynamic similitude

✓ Implications for laboratory experiments:

$$\Pi_{1m} = \varphi (\Pi_{2m}, \Pi_{3m}, \dots, \Pi_{km}) = \varphi (\Pi_{2p}, \Pi_{3p}, \dots, \Pi_{kp}) = \Pi_{1p}$$

✓ Limited or incomplete similitude. Necessity

3.3. Absolute and limited similitude. Froude and Reynolds incompatibility

✓ Incompatibility between Froude and Reynolds

- Model construction according to Froude

$$\boxed{Fr_m = Fr_p} \quad \boxed{\frac{U_m}{U_p} = \sqrt{\frac{l_m}{l_p}} = \lambda^{1/2}}$$

- Model construction according to Reynolds

$$\boxed{Re_m = Re_p} \quad \boxed{\frac{v_m}{v_p} = \frac{U_m}{U_p} \frac{l_m}{l_p}}$$

$$\boxed{\frac{v_m}{v_p} = \sqrt{\frac{l_m}{l_p}} \frac{l_m}{l_p} = \left(\frac{l_m}{l_p}\right)^{3/2} = \lambda^{3/2}}$$

3.3. Absolute and limited similitude. Froude and Reynolds incompatibility

✓ Types of limited similitude

- **Mach**: in any case where the compressibility of the fluid is significant (gases with high velocity)
- **Froude**: in any case where the gravity forces are important (dams, channels, buoyancy, open flows)
- **Reynolds**: in any case where the viscous effect is preponderant (fluid flow in closed ducts, hydraulic machines, aerodynamic tunnels with low Mach No)
- **Euler**: in any case where the fluid pressure state is important (cavitation phenomenon)

3.4. Scale distortion

- ✓ Necessity of scale distortion
- ✓ Example: simulation of the flow of a river

$$\lambda_1 = \frac{\text{length}_m}{\text{length}_p} = \frac{\text{width}_m}{\text{width}_p}$$

$$\lambda_2 = \frac{h_m}{h_p}$$

- **Model according to Froude:**

$$\frac{U_m}{\sqrt{g_m h_m}} = \frac{U_p}{\sqrt{g_p h_p}}$$

$$\frac{U_m}{U_p} = \sqrt{\frac{h_m}{h_p}}$$

Velocity ratio

$$\frac{Q_m}{Q_p} = \frac{A_m U_m}{A_p U_p} = \frac{\text{width}_m \times h_m}{\text{width}_p \times h_p} \sqrt{\frac{h_m}{h_p}} = \lambda_1 \lambda_2 \lambda_2^{1/2} = \lambda_1 \lambda_2^{3/2}$$

Flow rate ratio

3.5. Applications: external flow

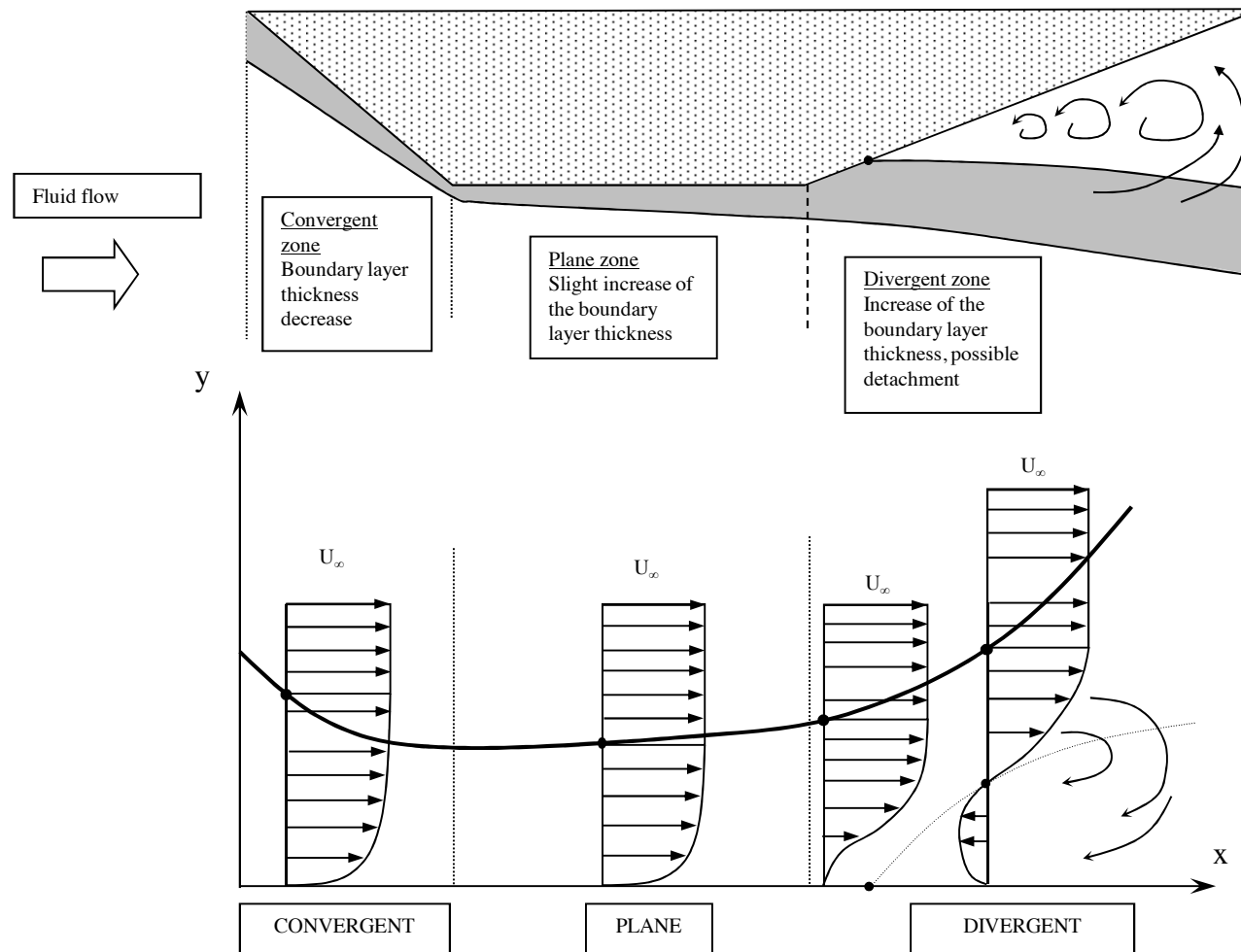
3.5.1. Drag force

3.5.2. Lift force

3.5.3. Flow induced vibration. Von Karman vortex street

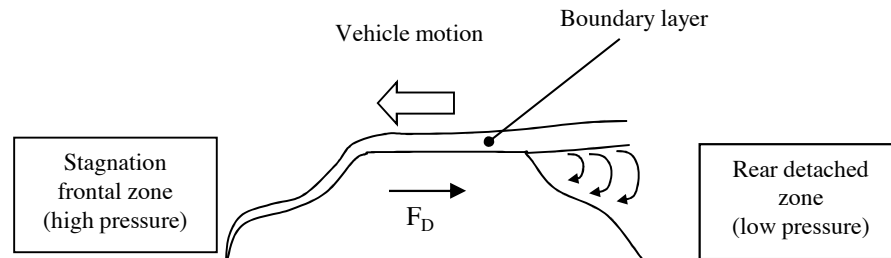
3.5.1. Drag force: friction and pressure

- ✓ Pressure drag or form drag: separation/detachment of the boundary layer



3.5.1. Drag force: friction and pressure

- ✓ Pressure drag or form drag: separation/detachment of the boundary layer



- ✓ General expression of the drag force

$$F_d = \frac{1}{2} C_d \rho A U^2$$

- Area:

1. Frontal area (projection on to the plane that is perpendicular to the motion direction) is used: in most cases, with totally immersed bodies.
2. Planform area (Proj. on to the horizontal plane) is used: with streamlined bodies (such as airfoils)
3. Wetted area (skin): with buoyant bodies (ships, boats)

3.5.1. Drag force: friction and pressure

- ✓ Friction and pressure:

$$F_d = F_{d,f} + F_{d,p}$$

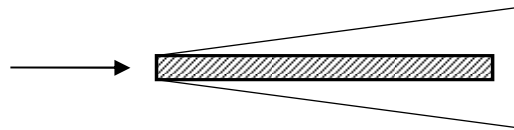
$$C_{d,f} = \frac{F_{d,f}}{\frac{1}{2}\rho U^2 A} \quad C_{d,p} = \frac{F_{d,p}}{\frac{1}{2}\rho U^2 A} \quad \Rightarrow \quad C_d = C_{d,f} + C_{d,p}$$

- ✓ Partially submerged bodies: “ wave drag”

$$F_w = \frac{1}{2} C_{d,w} \rho A U^2$$

3.5.1. Drag force: friction and pressure

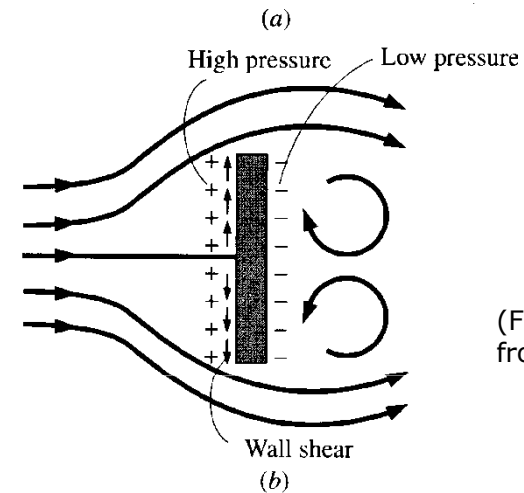
- ✓ Relative importance of pressure and friction drag:



Plane plate parallel to the flow: only friction

$$F_d = \iint_A \tau_p dA$$

“Friction drag”



(Figure taken from Cengel)

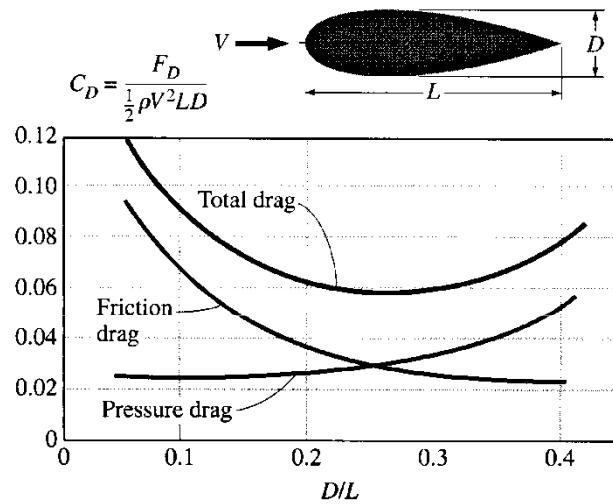
Flat plate normal to the flow: only pressure

$$F_d = \iint_{A_1} p dA - \iint_{A_2} p dA$$

“Pressure drag”

3.5.1. Drag force: friction and pressure

✓ Drag coefficient: form influence



(Figures taken from Cengel)

Representative drag coefficients C_D for various three-dimensional bodies based on the frontal area for $Re > 10^4$ unless stated otherwise (for use in the drag force relation $F_D = C_D A \rho V^2 / 2$ where V is the upstream velocity)

<p>Cube, $A = D^2$</p> <p>$C_D = 1.05$</p>	<p>Thin circular disk, $A = \pi D^2/4$</p> <p>$C_D = 1.1$</p>	<p>Cone (for $\theta = 30^\circ$), $A = \pi D^2/4$</p> <p>$C_D = 0.5$</p>																										
<p>Sphere, $A = \pi D^2/4$</p> <p>Laminar: $Re \leq 2 \times 10^5$ $C_D = 0.5$ Turbulent: $Re \geq 2 \times 10^6$ $C_D = 0.2$</p> <p>See Fig. 11-36 for C_D vs. Re for smooth and rough spheres.</p>	<p>Ellipsoid, $A = \pi D^2/4$</p> <table border="1"> <thead> <tr> <th rowspan="2">L/D</th> <th colspan="2">C_D</th> </tr> <tr> <th>Laminar $Re \leq 2 \times 10^5$</th> <th>Turbulent $Re \geq 2 \times 10^6$</th> </tr> </thead> <tbody> <tr><td>0.75</td><td>0.5</td><td>0.2</td></tr> <tr><td>1</td><td>0.5</td><td>0.2</td></tr> <tr><td>2</td><td>0.3</td><td>0.1</td></tr> <tr><td>4</td><td>0.3</td><td>0.1</td></tr> <tr><td>8</td><td>0.2</td><td>0.1</td></tr> </tbody> </table>	L/D	C_D		Laminar $Re \leq 2 \times 10^5$	Turbulent $Re \geq 2 \times 10^6$	0.75	0.5	0.2	1	0.5	0.2	2	0.3	0.1	4	0.3	0.1	8	0.2	0.1							
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<p>Hemisphere, $A = \pi D^2/4$</p> <p>$C_D = 0.4$</p> <p>$C_D = 1.2$</p>	<p>Finite cylinder, vertical, $A = LD$</p> <table border="1"> <thead> <tr> <th>L/D</th> <th>C_D</th> </tr> </thead> <tbody> <tr><td>1</td><td>0.6</td></tr> <tr><td>2</td><td>0.7</td></tr> <tr><td>5</td><td>0.8</td></tr> <tr><td>10</td><td>0.9</td></tr> <tr><td>40</td><td>1.0</td></tr> <tr><td>∞</td><td>1.2</td></tr> </tbody> </table> <p>Values are for laminar flow ($Re \leq 2 \times 10^5$)</p>	L/D	C_D	1	0.6	2	0.7	5	0.8	10	0.9	40	1.0	∞	1.2	<p>Finite cylinder, horizontal, $A = \pi D^2/4$</p> <table border="1"> <thead> <tr> <th>L/D</th> <th>C_D</th> </tr> </thead> <tbody> <tr><td>0.5</td><td>1.1</td></tr> <tr><td>1</td><td>0.9</td></tr> <tr><td>2</td><td>0.9</td></tr> <tr><td>4</td><td>0.9</td></tr> <tr><td>8</td><td>1.0</td></tr> </tbody> </table>	L/D	C_D	0.5	1.1	1	0.9	2	0.9	4	0.9	8	1.0
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<p>Streamlined body, $A = \pi D^2/4$</p> <p>$C_D = 0.04$</p>	<p>Parachute, $A = \pi D^2/4$</p> <p>$C_D = 1.3$</p>	<p>Tree, $A =$ frontal area</p> <table border="1"> <thead> <tr> <th>V, m/s</th> <th>C_D</th> </tr> </thead> <tbody> <tr><td>10</td><td>0.4-1.2</td></tr> <tr><td>20</td><td>0.3-1.0</td></tr> <tr><td>30</td><td>0.2-0.7</td></tr> </tbody> </table>	V , m/s	C_D	10	0.4-1.2	20	0.3-1.0	30	0.2-0.7																		
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<p>Rectangular plate, $A = LD$</p> <p>$C_D = 1.10 + 0.02 (L/D + D/L)$ for $1/30 < (L/D) < 30$</p>																												

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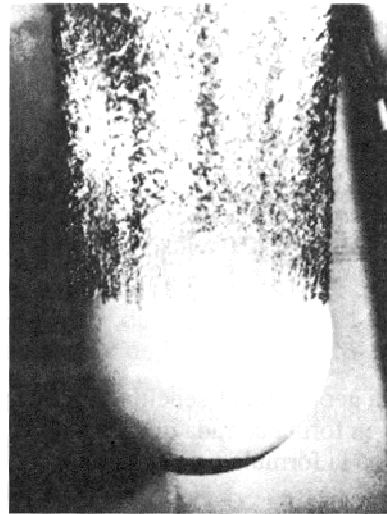
3.5.1. Flow over cylinders and spheres

✓ Type of boundary layer, detachment:



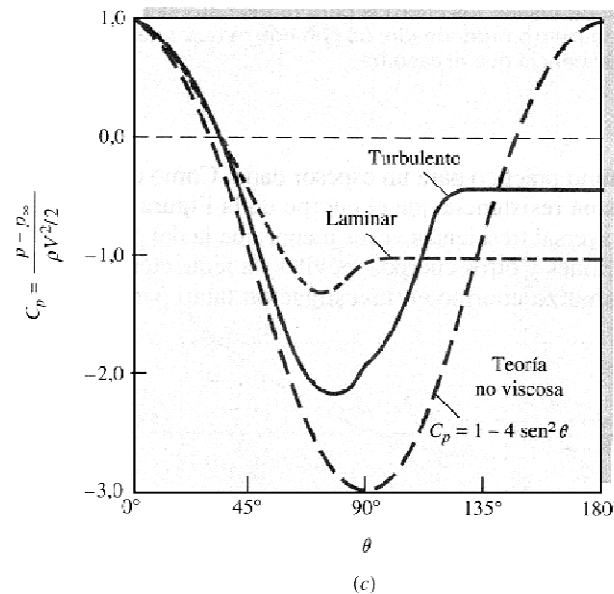
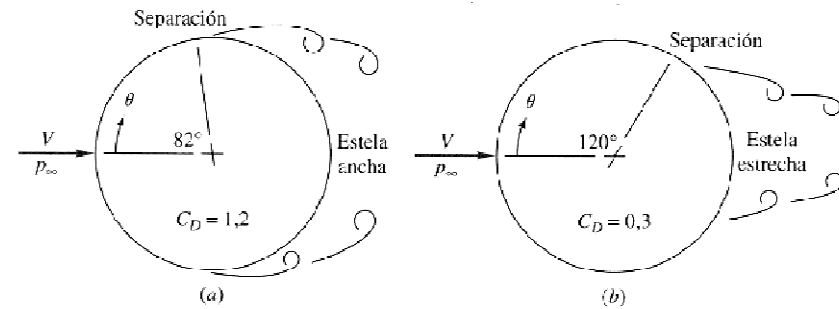
(a)

LAMINAR



(b)

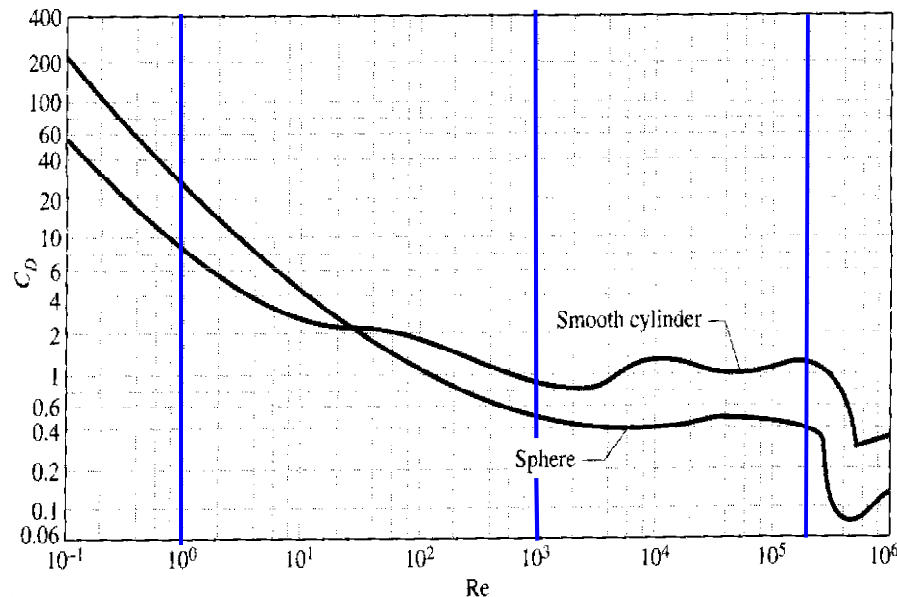
TURBULENT



(Figures taken from White)

3.5.1. Flow over cylinders and spheres

✓ Reynolds number, type of regime



(Figure taken from Cengel)

✓ Regions:

$Re < 1$: Dominant force: friction drag F_f .
It can be defined as: $C_d = k/Re$, in case of a sphere
 $C_d = 24/Re$

$$F = 6\pi R\mu U \quad (\text{"Stokes formula"})$$

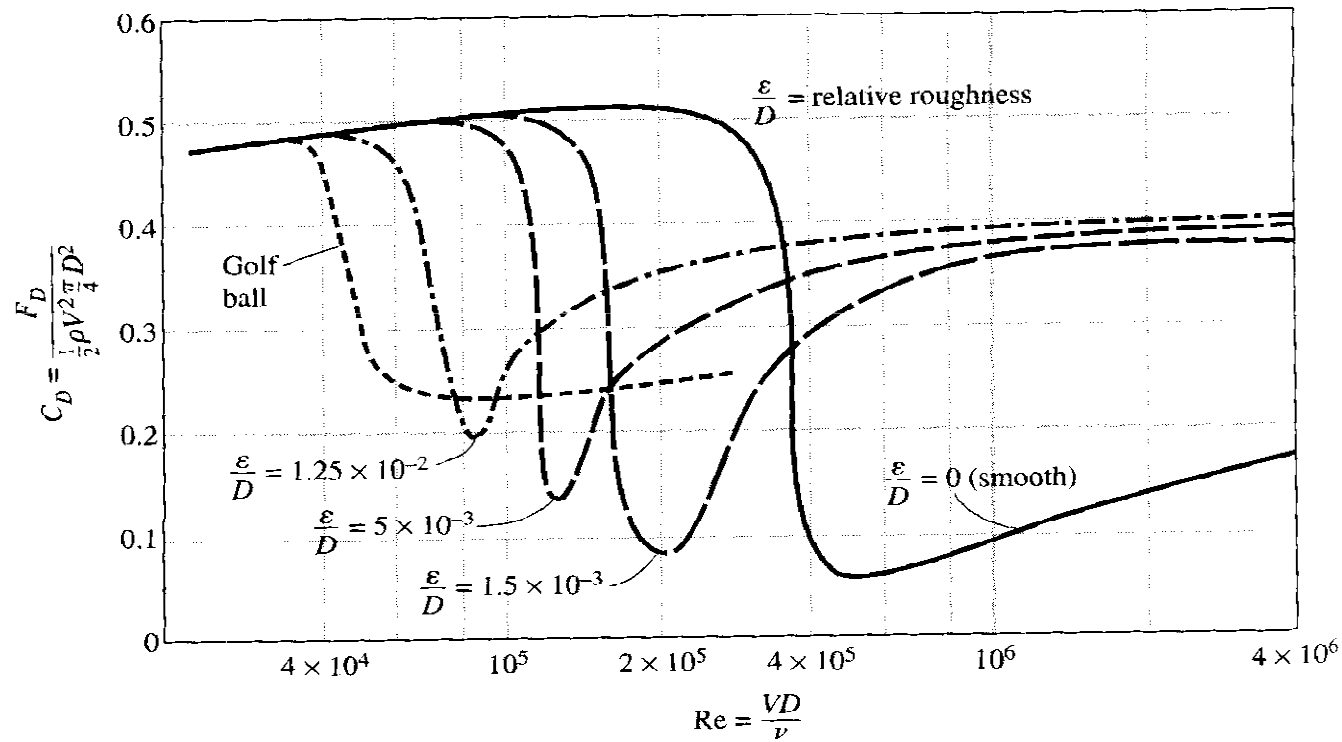
$1 < Re < 1000$: Transition zone, friction drag F_f and pressure drag F_p simultaneously intervening.

$1000 < Re < 2 \times 10^5$: Dominant force: pressure drag F_p

$Re > 2 \times 10^5$: The boundary layer becomes turbulent

3.5.1. Flow over cylinders and spheres

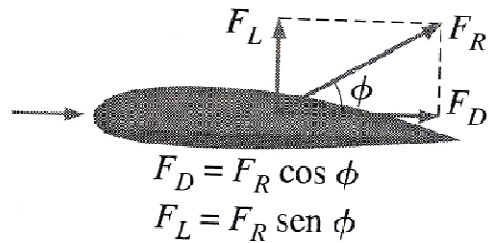
- ✓ Effect of roughness



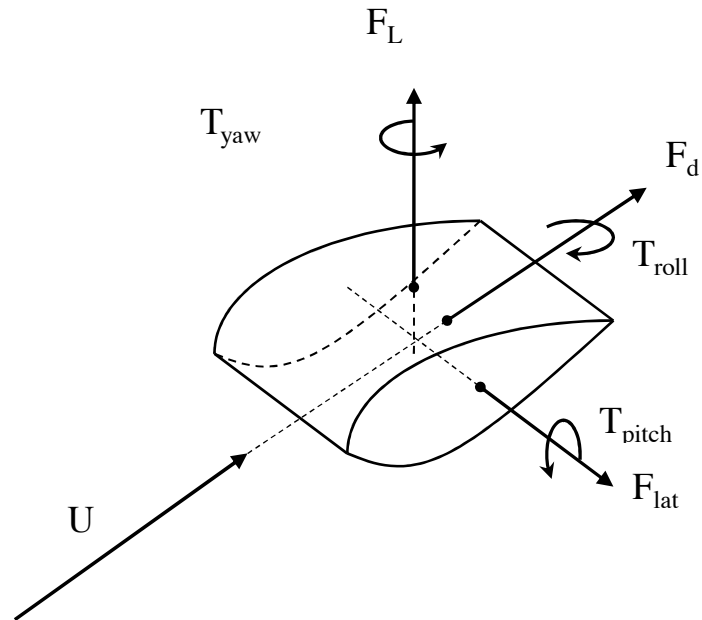
(Taken from Cengel)

3.5.2. Lift force in airfoils (wings)

✓ Airfoil:



(Figure taken from Cengel)

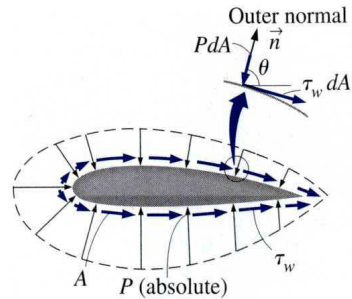


F_d : Drag force (parallel to flow)
 F_L : Lift force (normal to flow)
 F_{lat} : Lateral force

Roll torque: around flow direction
Yaw torque: around lift direction
Pitch torque: around lateral force direction

3.5.2. Lift force in airfoils (wings)

✓ Source:



(Figure taken from Cengel)

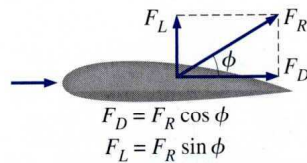
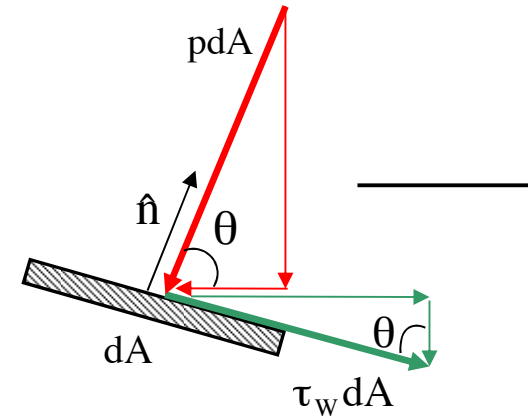


FIGURE 11-5

The pressure and viscous forces acting on a two-dimensional body and the resultant lift and drag forces.

- Differential area element "dA":



Shear force: $\tau_w dA$

Normal force: $p dA$

- Total forces:

$$F_d = \iint_A (-p \cos \theta + \tau_w \sin \theta) dA$$

$$F_L = \iint_A (-p \sin \theta - \tau_w \cos \theta) dA$$

$$\left\{ \begin{array}{l} \text{Drag contribution:} \\ dF_d = -p dA \cos \theta + \tau_w dA \sin \theta \\ \text{Lift contribution:} \\ dF_L = -p dA \sin \theta - \tau_w dA \cos \theta \end{array} \right.$$

3.5.2. Lift force in airfoils (wings)

✓ Calculation:

$$F_L = \frac{1}{2} C_L \rho A U^2$$

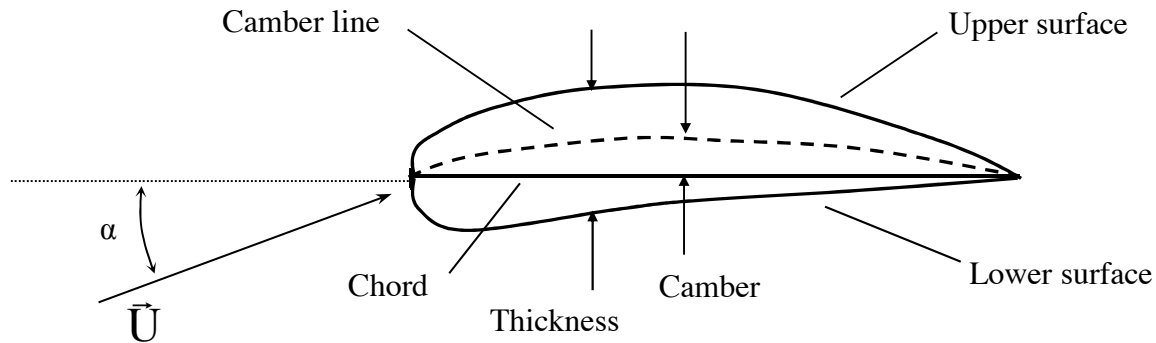
C_L : Lift coefficient (for a specific body as a function of the angle of attack α and Reynolds number (less important))

ρ : fluid density

A: planform area

U: freestream velocity

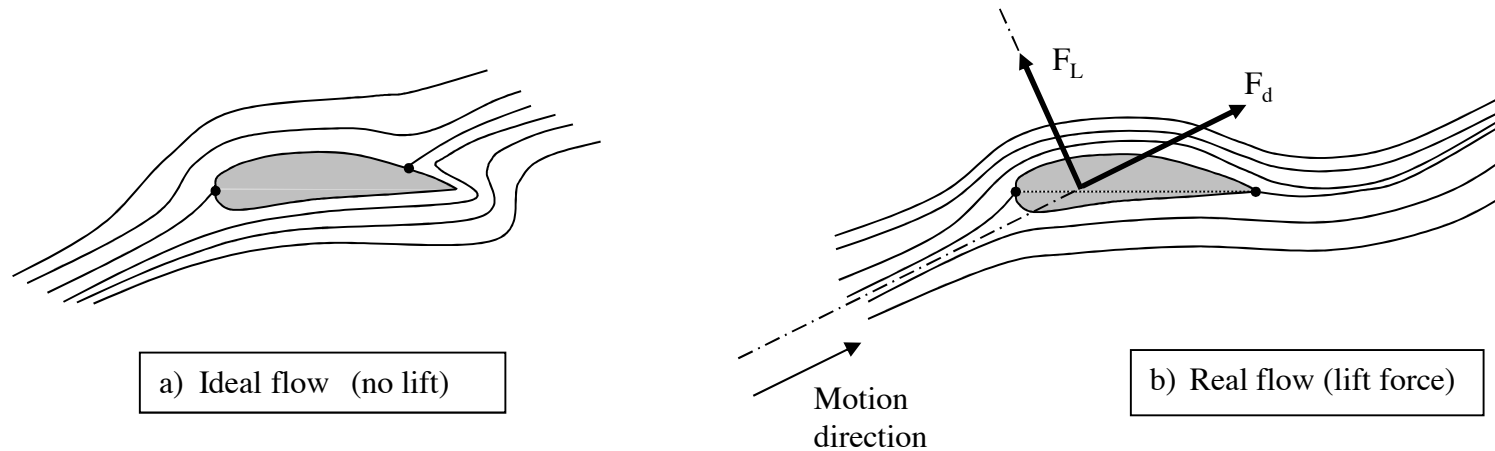
✓ Airfoil terminology:



- Chord
- Camber line
- Camber
- Angle of attack " α "
- Wingspan (or span)
- Leading edge
- Trailing edge

3.5.2. Lift force in airfoils (wings)

✓ Source:



Lift force increase:

- (1) Curvature of the airfoil (longer length of the upper surface in comparison to the lower surface)
- (2) Angle of attack (up to a value of 15° aprox.)
- (3) Use of movable leading edge and trailing edge flaps

3.5.2. Lift force in airfoils (wings)

✓ Source:

Lift force increase:

(3) Use of flaps

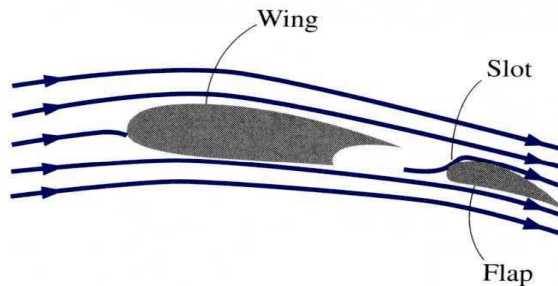
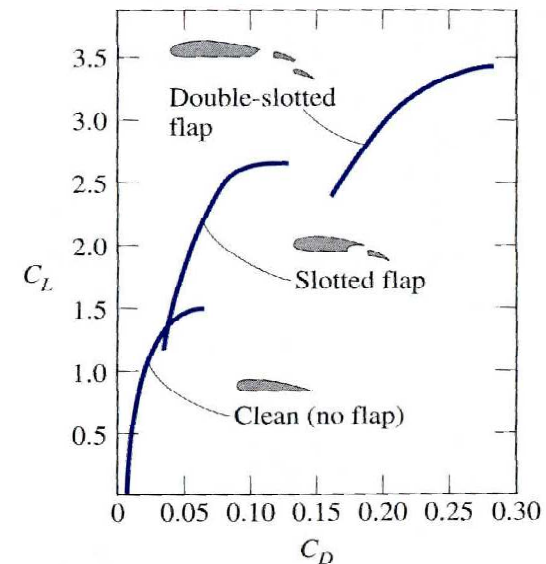
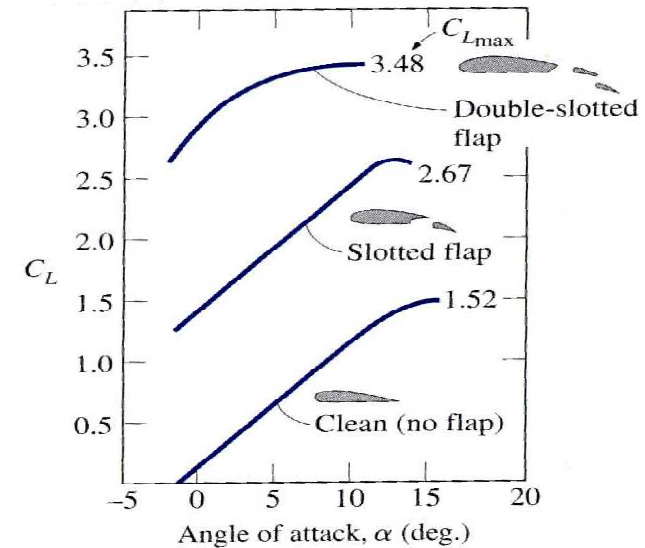


FIGURE 11-46

A flapped airfoil with a slot to prevent the separation of the boundary layer from the upper surface and to increase the lift coefficient.

(Figures taken from Cengel)

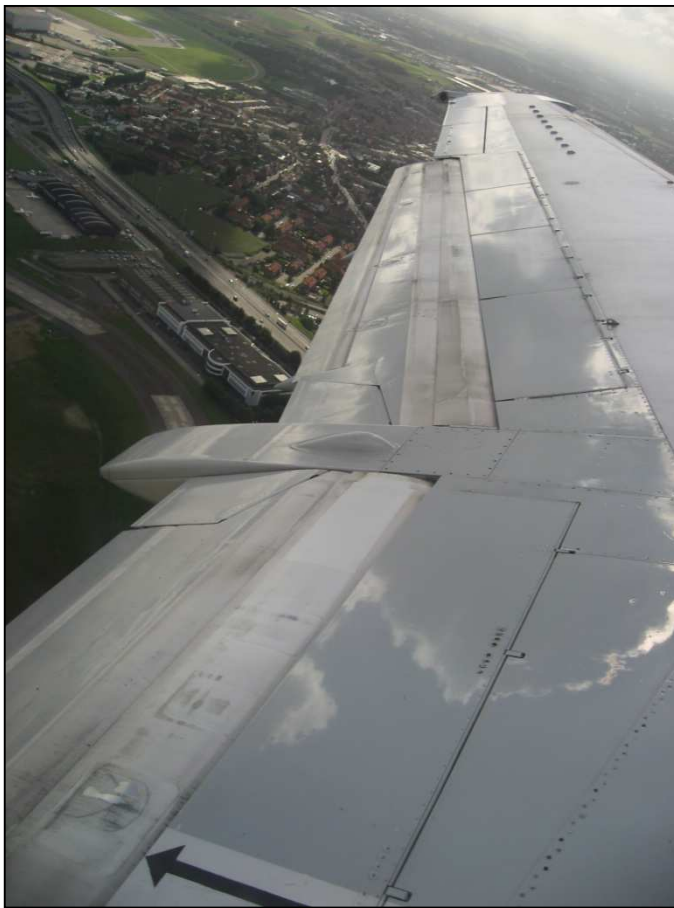


3.5.2. Lift force in airfoils (wings)

✓ Source:

Lift force increase:

(3) Use of flaps



(Photographs G. A. Esteban)

3.5.2. Lift by Magnus effect

- ✓ Cylinder rotation: lift generation

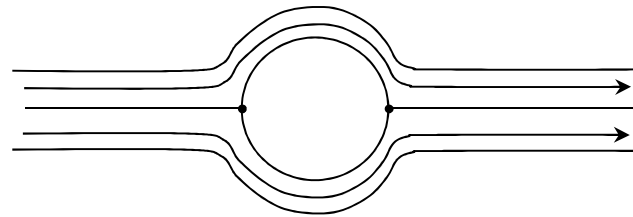


Figure 8.9 Fluid streamlines around a static cylinder

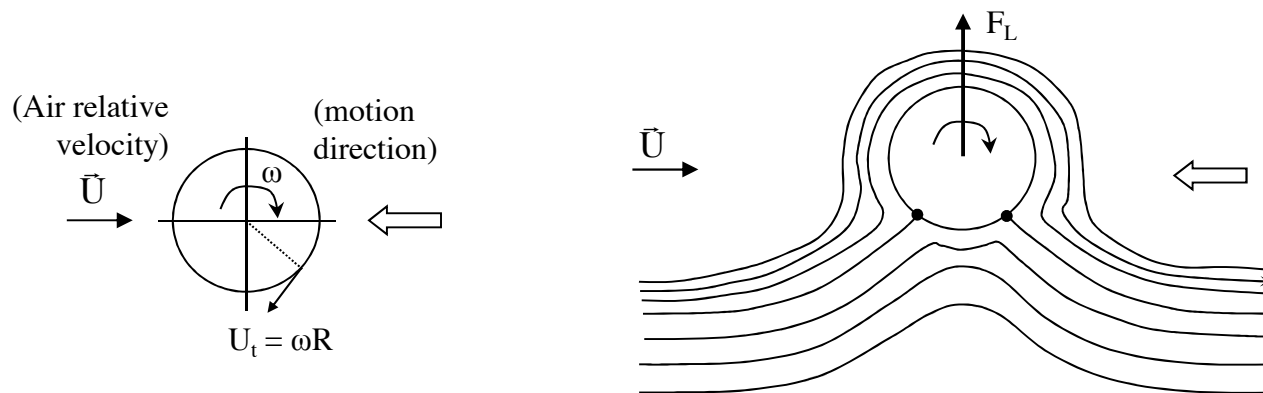
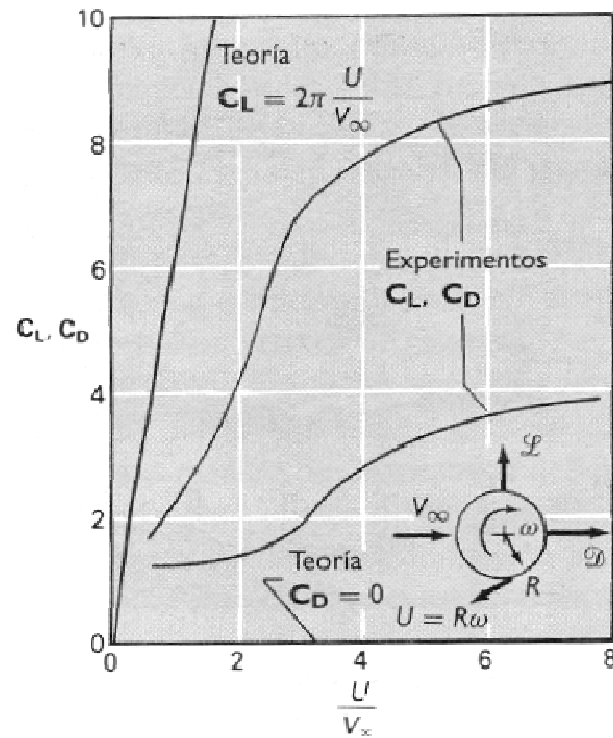


Figure 8.10 Stagnation points modification because of the spinning

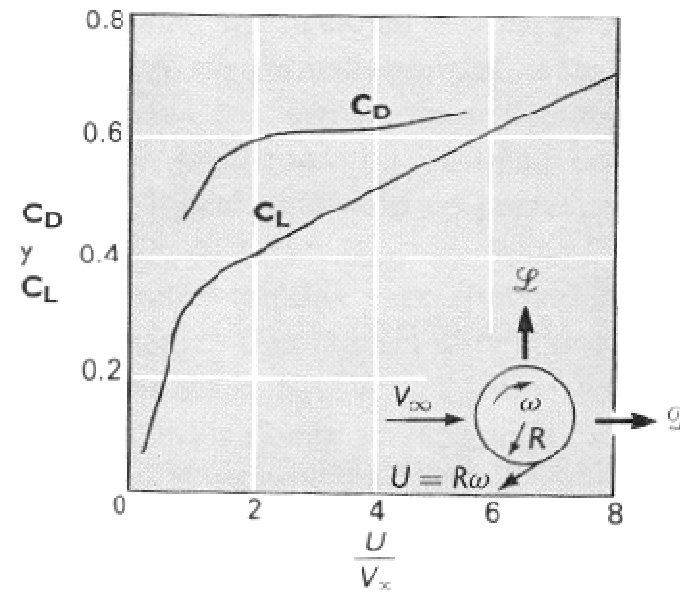
3.5.2. Lift by Magnus effect

- ✓ Drag and lift coefficients (experimental)



SMOOTH CYLINDER

(Figures taken from Gerhart)



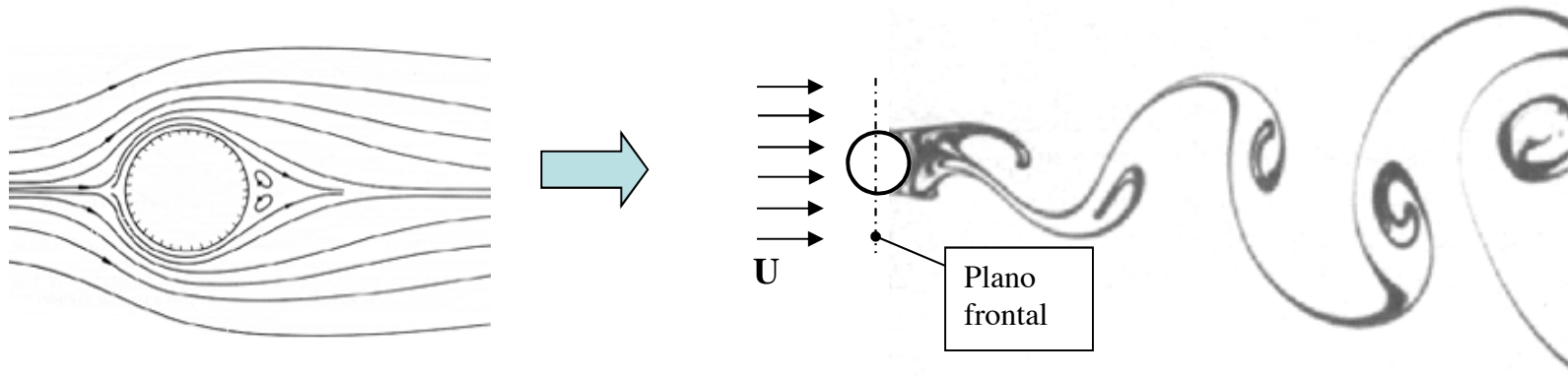
SMOOTH SPHERES

- Theoretical value (cylinder):

$$C_L = \frac{2\pi\omega R}{U_\infty}$$

3.5.3. Von-Karman vortex street

✓ Phenomenon



✓ Dimensional analysis

The vortex emission rate "n" (Hz) o (s^{-1}) depends on free stream flow velocity "U", the cylinder diameter "D", fluid viscosity " μ ", fluid density " ρ ".

Π groups describing the phenomenon:

Strouhal number: $S = \frac{nD}{U}$

Reynolds number: $Re = \frac{UD\rho}{\mu}$

3.5.3. Von-Karman vortex street

✓ Relationship between dimensionless numbers:

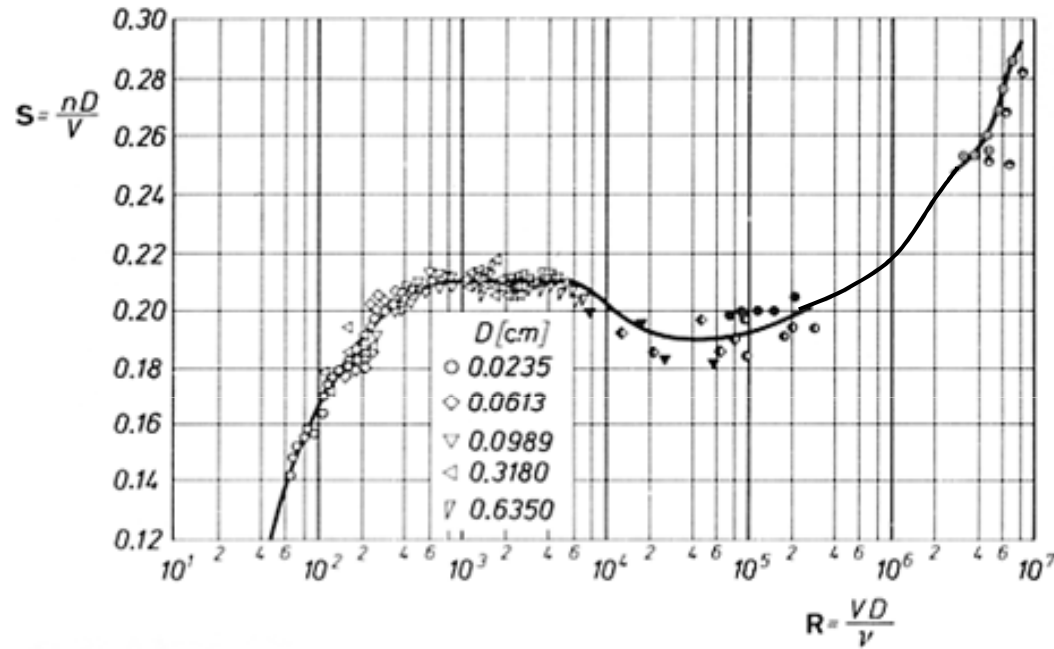


Figure 8.13 Dependence of the vortex emission rate

Figure 8.14 Takoma bridge (Washington) fallen down by "weak" wind 7th of November, 1940.

