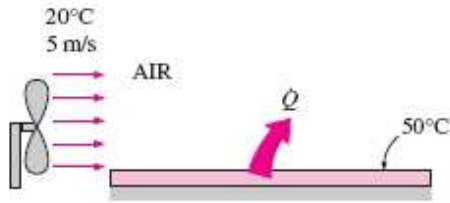


6. GAIA

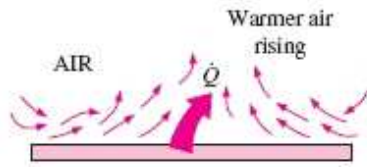
KONBEKZIOAREN OINARRIAK

6.0 - HELBURUAK

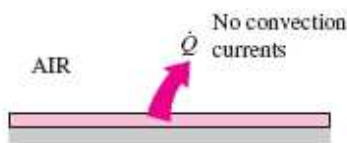
- **Konbekzioaren** mekanismo fisikoak eta sailkapena ulertu.
- Gainazalen gaineko fluxuen **abiaduraren mugalde-geruza** eta **mugalde-geruza termikoa** irudikatu.
- Reynoldsen, Prandtlen eta Nusselten **zenbaki dimentsiogabeen** ezagutza praktikoa izan.
- Fluxu **laminarrak** eta **turbulentuak** bereizi, eta fluxu turbulentuen momentu- eta bero-transferentziako mekanismoak ulertu
- Konbekzioa deskribatzen duten **ekuazio diferentzialak garatu**, masa-, momentu- eta energia-balantzeetan oinarrituta
- Konbekzio-ekuazioak **dimentsiogabetu**, eta marruskaduraren eta bero-transferentziaren koefizienteen forma **funtzionalak** lortu.
- Momentu- eta bero-transferentziaren arteko **analogiak** erabili, eta bero-transferentziaren koefizientea kalkulatu, marruskadura-koefizientetik abiatuta



(a) Forced convection



(b) Free convection



(c) Conduction

Konbekzio bero-transferentzia hurrengo parametroen araberakoa da:

Biskositate dinamikoa μ_{fluid}

Eroankortasun termikoa k_{fluid}

Dentsitatea ρ_{fluid}

Bero espezifikoa $C_{p,fluid}$

Jariakinaren abiadura V_{fluid}

Konfigurazio geometrikoa

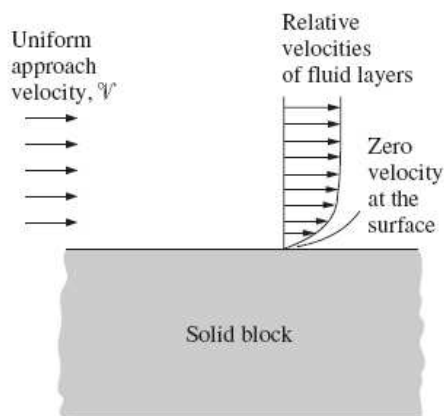
Solidoaren zimurtasuna

Jariakinaren fluxu mota (laminarra edo turbulentua)

Newtonen hozte legea

$$\dot{Q} = h \cdot A \cdot (T_s - T_\infty) \quad [W]$$

TERMOTEKNIA



Irristadurarik ezaren baldintza

Mugalde geruza

$$\dot{q}_{cond} = -k_{fluid} \cdot \left. \frac{\partial T}{\partial y} \right|_{y=0} \quad [W/m^2]$$

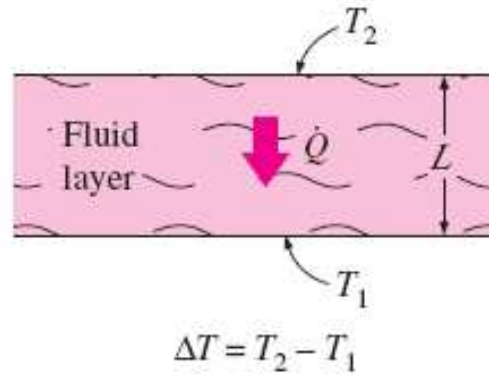
$$\dot{q}_{conv} = h \cdot (T_s - T_\infty) \quad [W/m^2]$$

$$h = \frac{-k_{fluid} \cdot (\partial T / \partial y)_{y=0}}{T_s - T_\infty} \quad [W/m^2 \cdot ^\circ C]$$

Bataz besteko konbekzio koefizientea – Konbekzio koefiziente lokala

NUSSELTEN ZENBAKIA

$$Nu = \frac{h \cdot L_c}{k}$$



$$\left. \begin{aligned} \dot{q}_{conv} &= h \cdot \Delta T \\ \dot{q}_{cond} &= k \cdot \frac{\Delta T}{L} \end{aligned} \right\} \Rightarrow \frac{\dot{q}_{conv}}{\dot{q}_{cond}} = \frac{h \cdot \Delta T}{k \cdot \Delta T / L} = \frac{h \cdot L}{k} = Nu$$

6.2 – FLUIDO-FLUXUEN SAILKAPENA

Fluxu likatsua vs ez-likatsua.

Barne-fluxua vs kanpo-fluxua.

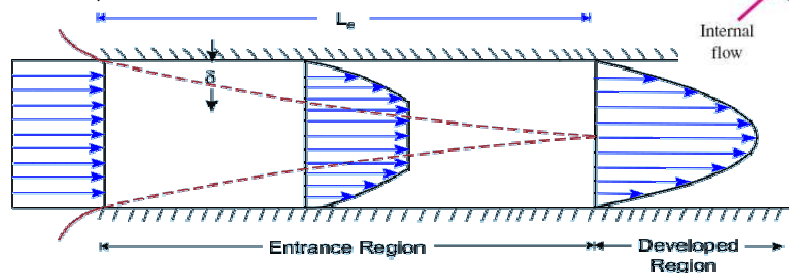
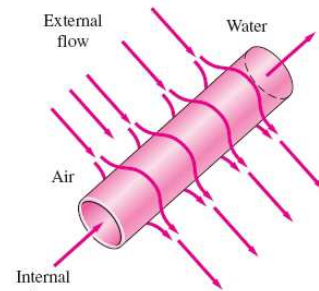
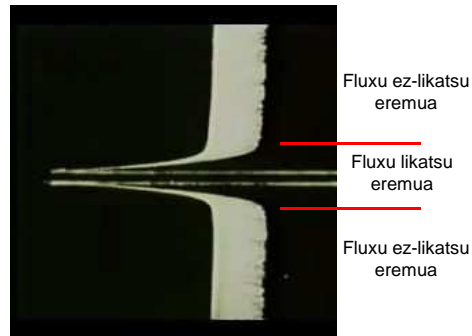
Fluxu konprimagarria vs konprimaezina.

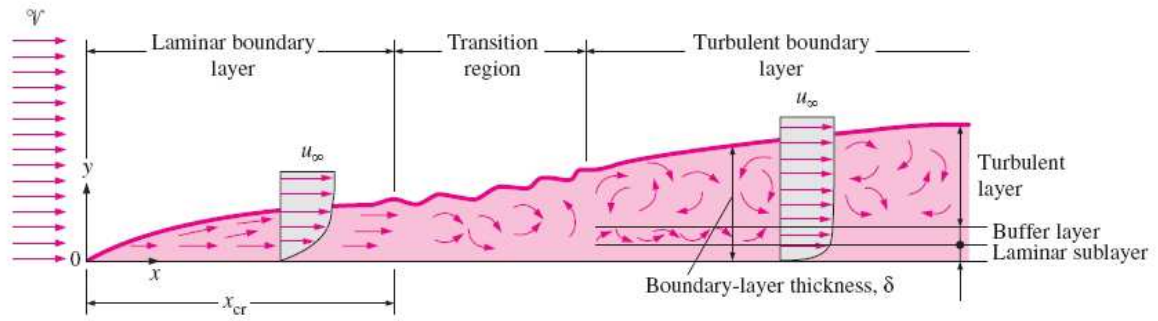
Fluxu laminarra vs turbulentua.

Fluxu naturala vs behartua.

Fluxu geldikorra vs ez-geldikorra (iragankorra).

Dimensio bakarreko, biko eta hiruko fluxuak.



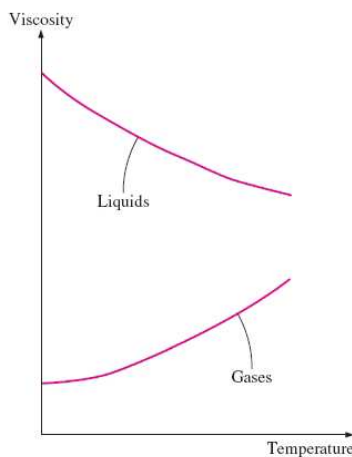


GAINAZALEKO EBAKIDURA-TENTSIOA $\tau = \mu \frac{\partial u}{\partial y} \Big|_{y=0}$ [N/m²] Jariakin newtondarrak

Biskositate dinamikoa: μ [kg/m·s] [poise] 1 poise = 0,1 kg /m·s

Biskositate zinematikoa: $\nu = \frac{\mu}{\rho}$ [m²/s] [stoke] 1 stoke = 10⁻⁴ m²/s

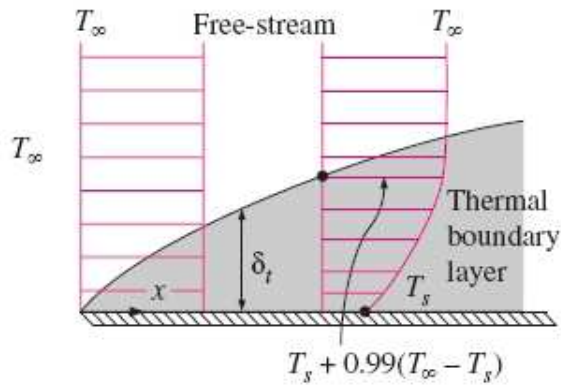
GAINAZALEKO EBAKIDURA-TENTSIOA



Fluid	Dynamic viscosity μ , kg/m · s
Glycerin:	
-20°C	134.0
0°C	12.1
20°C	1.49
40°C	0.27
Engine oil:	
SAE 10W	0.10
SAE 10W30	0.17
SAE 30	0.29
SAE 50	0.86
Mercury	0.0015
Ethyl alcohol	0.0012
Water:	
0°C	0.0018
20°C	0.0010
100°C (liquid)	0.0003
100°C (vapor)	0.000013
Blood, 37°C	0.0004
Gasoline	0.00029
Ammonia	0.00022
Air	0.000018
Hydrogen, 0°C	0.000009

$\tau_s = C_f \frac{\rho \cdot V^2}{2}$

Marrusadura-koefizientea



Typical ranges of Prandtl numbers for common fluids

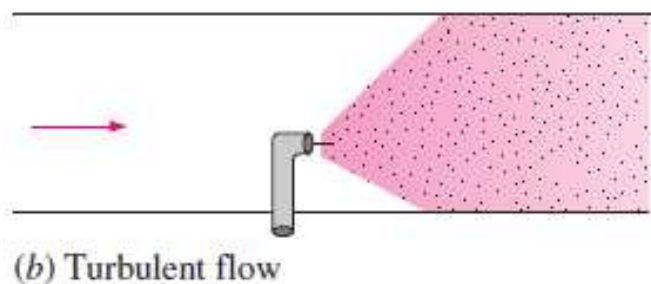
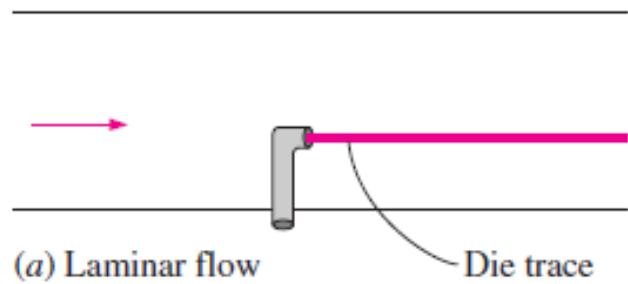
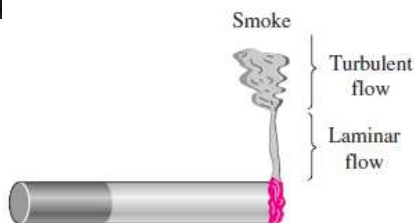
Fluid	Pr
Liquid metals	0.004–0.030
Gases	0.7–1.0
Water	1.7–13.7
Light organic fluids	5–50
Oils	50–100,000
Glycerin	2000–100,000

PRANDTLEN ZENBAKIA

$$Pr = \frac{\nu}{\alpha} = \frac{\mu \cdot c_p}{k}$$

$$Pr = \frac{\text{Momentuaren difusibitate molekularra}}{\text{Beroaren difusibitate molekularra}}$$

6.5 – FLUXU LAMINARRAK ETA TURBULENTOAK

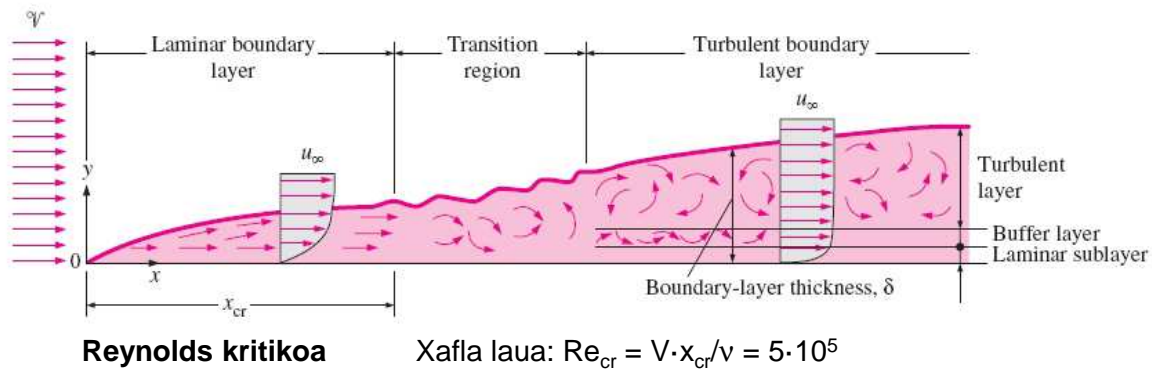


REYNOLDSSEN ZENBAKIA

$$Re = \frac{V \cdot L_c}{\nu} = \frac{\rho \cdot V \cdot L_c}{\mu}$$

$Re = \frac{\text{Inertzi indarrak}}{\text{Biskosite - indarrak}}$

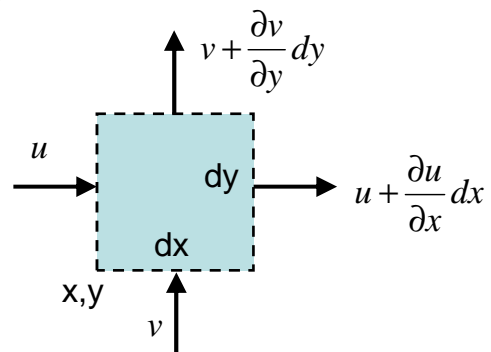
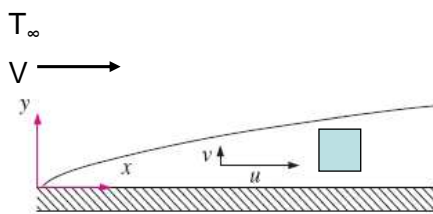
Re baxuak → Fluxu laminarra
 Re altuak → Fluxu turbulenta



6.7 – KONBEKZIO-EKUAZIO DIFERENTZIALEN GARAPENA

Hipotesiak: Fluxu geldikorra bi dimentsiotan
 Jariakin Newtondarra
 Propietateak konstante (ρ, ν, k, \dots)

JARRAITUTASUN-EKUAZIOA

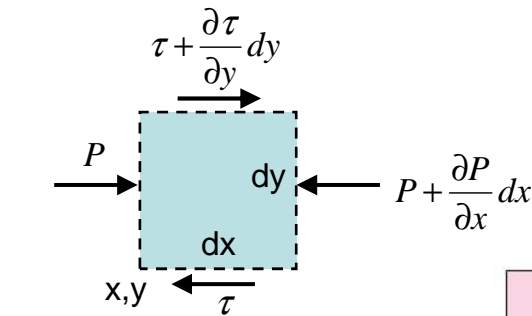


Masa-emia kontrol-bolumenaren barrurantz = Masa-emia kontrol-bolumenaren kanporantz

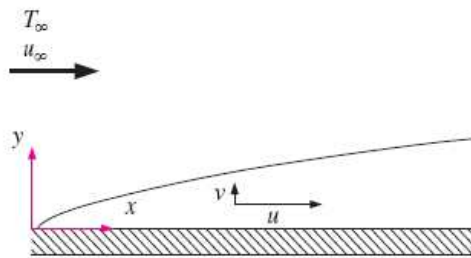
$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

MOMENTU-EKUAZIOAK

$$\left(\text{Masa} \right) \left(\begin{array}{c} \text{Azelerazioa, zehaztutako} \\ \text{noranzkoan} \end{array} \right) = \left(\begin{array}{c} \text{Noranzko horretan eragiten duen indar} \\ \text{garbia (gorputzekoa eta gainazalekoa)} \end{array} \right)$$



$$\rho \left(u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} \right) = \mu \frac{\partial^2 u}{\partial y^2} - \frac{\partial P}{\partial x}$$



- 1) Velocity components:
 $v \ll u$
- 2) Velocity gradients:
 $\frac{\partial v}{\partial x} \approx 0, \frac{\partial v}{\partial y} \approx 0$
 $\frac{\partial u}{\partial x} \ll \frac{\partial u}{\partial y}$
- 3) Temperature gradients:
 $\frac{\partial T}{\partial x} \ll \frac{\partial T}{\partial y}$

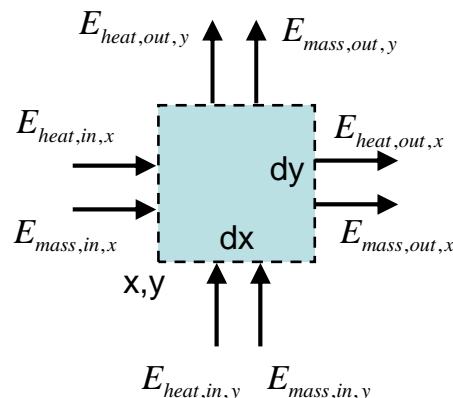
$$\frac{\partial P}{\partial y} = 0$$

ENERGIA-KONTSERBAZIOAREN EKUAZIOA

$$E_{in} - E_{out} = \Delta E_{sist} \xrightarrow{\text{Egoera egonkorrean:}} \dot{E}_{in} - \dot{E}_{out} = 0$$

$$\left(\dot{E}_{in} - \dot{E}_{out} \right)_{\text{by heat}} + \left(\dot{E}_{in} - \dot{E}_{out} \right)_{\text{by work}} + \left(\dot{E}_{in} - \dot{E}_{out} \right)_{\text{by mass}} = 0$$

$$e_{stream} = h + e_c + e_p = h + \frac{1}{2} V^2 + g \cdot z \approx h = c_p \cdot T$$



ENERGIA-KONTSERBAZIOAREN EKUAZIOA

$$(\dot{E}_{in} - \dot{E}_{out})_{\text{by mass}} = -\rho \cdot c_p \cdot \left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) dx \cdot dy$$

$$(\dot{E}_{in} - \dot{E}_{out})_{\text{by heat}} = k \cdot \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) dx \cdot dy$$

$$(\dot{E}_{in} - \dot{E}_{out})_{\text{by work}} = \mu \cdot \Phi \cdot dx \cdot dy = \mu \cdot \left(\frac{\partial u}{\partial y} \right)^2 \cdot dx \cdot dy \quad \text{Abiadura baxuetan mesprezagarria}$$

$$\dot{E}_{in} - \dot{E}_{out} = 0 \quad \rightarrow \quad \rho \cdot c_p \cdot \left(u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} \right) = k \cdot \left(\frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right)$$

Mugalde geruzaren barnean oso txikia

6.9 – KONBEKZIO-EKUAZIO DIMENTSIOGABEAK ETA ANTZEKOTASUNA

Aldagai guztiak dimentsiogabetuz:

$$x^* = \frac{x}{L}; \quad y^* = \frac{y}{L}; \quad u^* = \frac{u}{V}; \quad v^* = \frac{v}{V}; \quad P^* = \frac{P}{\rho \cdot V^2}; \quad T^* = \frac{T - T_s}{T_\infty - T_s};$$

Jarraitutasuna:

$$\frac{\partial u^*}{\partial x^*} + \frac{\partial v^*}{\partial y^*} = 0$$

Momentua:

$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = \frac{1}{\text{Re}_L} \frac{\partial^2 u^*}{\partial y^{*2}} - \frac{\partial P^*}{\partial x^*}$$

Energia:

$$u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{1}{\text{Re}_L \cdot \text{Pr}} \left(\frac{\partial^2 T^*}{\partial y^{*2}} \right)$$

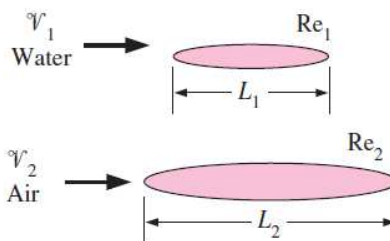
$$u^*(0, y^*) = 1; \quad u^*(x^*, 0) = 0; \quad u^*(x^*, \infty) = 1; \quad v^*(x^*, 0) = 0;$$

$$T^*(0, y^*) = 1; \quad T^*(x^*, 0) = 0; \quad T^*(x^*, \infty) = 1;$$

“Bi fenomeno fisiko *antzekoak* dira, baldin eta deskribatzen dituzten ekuazio diferentzialen eta mugalde-baldintzen forma dimentsiogabe berak badituzte.”

$$\begin{array}{l}
 L \\
 V \\
 T_{\infty} \\
 T_S \\
 \nu \\
 \alpha
 \end{array}
 \longrightarrow
 \begin{array}{l}
 Re_L \\
 Pr
 \end{array}$$

“Geometria jakin batean, antzekotasun-parametroen balio **bera** duten problemek **ebazpen berdin-berdinak** dituzte.”



Baldin $Re_1 = Re_2$ orduan $C_{f1} = C_{f2}$

6.10 – MARRUSKADURA- ETA KONBEKZIO-KOEFIZIENTEEN FORMA FUNTZIONALAK

3 ekuazio dimentsiogabeak	}	3 funtzio ezezagun:	$u^*, v^* \text{ y } T^*$
		2 aldagai independente:	x^*, y^*
		2 parametro:	$Re_L \text{ y } Pr$

Lehenengo bi ekuazioetatik: $u^* = f_1(x^*, y^*, Re_L)$

Ebakidura-tentsioa: $\tau_s = \mu \frac{\partial u}{\partial y} \Big|_{y=0} = \frac{\mu \cdot V}{L} \frac{\partial u^*}{\partial y^*} \Big|_{y^*=0} = \frac{\mu \cdot V}{L} f_2(x^*, Re_L)$

Marruskadura-koefizientea:

$$C_{f,x} = \frac{\tau_s}{\rho \cdot V^2 / 2} = \frac{\mu \cdot V / L}{\rho \cdot V^2 / 2} f_2(x^*, Re_L) = f_3(x^*, Re_L)$$

6.10 – MARRUSKADURA- ETA KONBEKZIO-KOEFIZIENTEEN FORMA FUNTZIONALAK

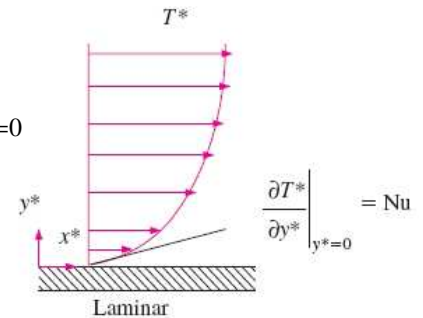
19/22

Azkenengo ekuaziotik.: $T^* = g_1(x^*, y^*, Re_L, Pr)$

Konbekzio koefizientea: $y = 0 \Rightarrow \dot{q}_{cond} = \dot{q}_{conv} \Rightarrow -k_{fluid} \cdot \frac{\partial T}{\partial y} \Big|_{y=0} = h \cdot (T_s - T_\infty)$

$$h = -\frac{k \frac{\partial T}{\partial y} \Big|_{y=0}}{T_s - T_\infty} = -\frac{k \frac{T_\infty - T_s}{L} \cdot \frac{\partial T^*}{\partial y^*} \Big|_{y^*=0}}{T_s - T_\infty} = \frac{k}{L} \frac{\partial T^*}{\partial y^*} \Big|_{y^*=0}$$

$$Nu_x = \frac{h \cdot L}{k} = \frac{\partial T^*}{\partial y^*} \Big|_{y^*=0} = g_2(x^*, Re_L, Pr)$$



0 eta 1 artean, X-ekiko integratuz

$$C_f = f_4(Re_L) \quad Nu = g_3(Re_L, Pr)$$

Sarritan:

$$Nu = C \cdot Re_L^m \cdot Pr^n$$

6.11 – MOMENTU- ETA BERO TRANSFERENTZIAREN ARTEKO ANALOGIAK

20/22

Helburua: C_f eta Nu lortu



Reynoldsen analogia

Chilton-Colburnen analogia

Baldin $Pr = 1$ eta $\frac{\partial P^*}{\partial x^*} = 0$

Momentua:

$$u^* \frac{\partial u^*}{\partial x^*} + v^* \frac{\partial u^*}{\partial y^*} = \frac{1}{Re_L} \frac{\partial^2 u^*}{\partial y^{*2}}$$

Energia:

$$u^* \frac{\partial T^*}{\partial x^*} + v^* \frac{\partial T^*}{\partial y^*} = \frac{1}{Re_L} \frac{\partial^2 T^*}{\partial y^{*2}}$$

Mugalde baldintzak:

$$u^*(0, y^*) = 1; \quad u^*(x^*, 0) = 0; \quad u^*(x^*, \infty) = 1;$$

$$T^*(0, y^*) = 1; \quad T^*(x^*, 0) = 0; \quad T^*(x^*, \infty) = 1;$$

u^* eta T^* berdinak behar dute izan

$$\frac{\partial u^*}{\partial y^*} \Big|_{y^*=0} = \frac{\partial T^*}{\partial y^*} \Big|_{y^*=0}$$

$$C_{f,x} \frac{Re_L}{2} = Nu_x \quad (Pr = 1)$$

Reynoldsen analogia

$$\frac{C_{f,x}}{2} = St_x \quad (Pr = 1)$$

$$St = \frac{h}{\rho \cdot c_p \cdot V} = \frac{Nu}{Re_L \cdot Pr}$$

Stantonen zenbakia

Prandtl-en zenbakia zuzenduz

Reynoldsen analogia eraldatua edo
Chilton-Colburnen analogia

$$C_{f,x} \frac{Re_L}{2} = Nu_x \cdot Pr^{-1/3}$$

$$\frac{C_{f,x}}{2} = \frac{h_x}{\rho \cdot c_p \cdot V} \cdot Pr^{2/3} \equiv j_H$$

Colburnen j faktorea

Baliogarria: $0,6 < Pr < 60$ tartean

Fluxu turbulenta presio-gradienteak badaude ere

Fluxu laminarra baldin eta $\partial P^* / \partial x^* = 0$

6.12 – IRAKATSIKO EZ DIREN ATALAK

6.6 atala – Fluxu turbulentuen bero- eta momentu-transferentzia

6.8 atala – Xafla lau baten konbekzio-ekuazioen ebazpenak